

Stochastic resonance

Motivational example:

Ice ages

• $x(t) \equiv$ global ice volume

$|x(\omega)|^2$: peak at $\Omega = \frac{2\pi}{T}$, $T \approx 10^5$ y.

• external driving?

eccentricity of earth orbit

oscillates with $T = 0.96 \cdot 10^5$ y,

induces $O(0.3\text{K})$ -temperature variation by change in irradiation

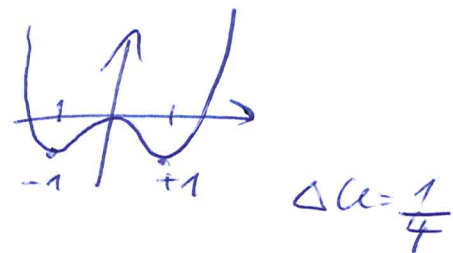
→ too weak to account for $\sim 10\text{K}$ temperature change during ice age.

• add noise!

Minimal model:

- diffusion in double-well potential with weak periodic driving.

$$U_0 = -\frac{1}{2}x^2 + \frac{1}{4}x^4$$



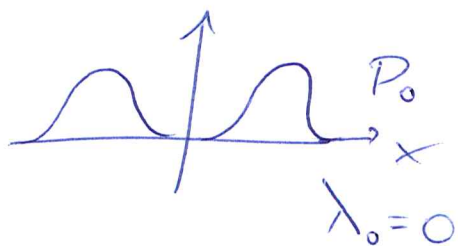
$$U = U_0 + A \cos \Omega t$$

$$\dot{x} = -\frac{\partial U}{\partial x} + \xi, \quad \langle \xi(t) \xi(t') \rangle = 2D \delta(t-t')$$

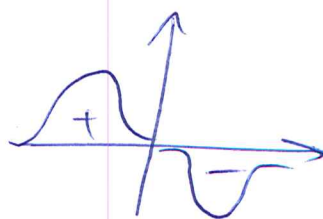
Kramer's rate

$$\begin{aligned} r_0 &= \frac{1}{2\pi} \sqrt{|U''(0) + U''(1)|} \exp\left(-\frac{\Delta U_0}{D}\right) \\ &= \frac{\sqrt{2}}{2\pi} \exp\left(-\frac{1}{4D}\right) \end{aligned}$$

Eigenfunctions of Fokker-Planck operator



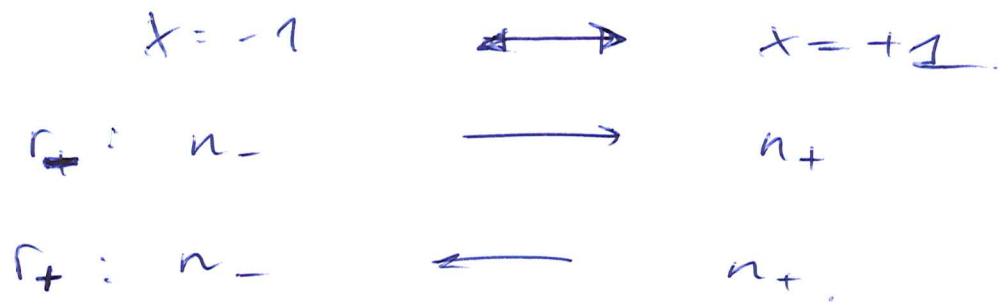
≡ steady-state distribution



≡ lowest mode

$$\lambda_1 = -r_0 < 0$$

two-state (telegraph) process.



$$\begin{aligned}
 \dot{n}_+ - \dot{n}_- &= r_- n_- - r_+ n_+ \\
 &= r_- - \underbrace{(r_+ + r_-)}_r n_+
 \end{aligned}$$

formal solution

$$n_+(t) = \left(\exp \int_{t_0}^t dt' r \right) n_+(t_0)$$

$$+ \exp \left(- \int_{t_0}^t dt' r \right) \cdot \int_{t_0}^t dt' r_-(t') \cdot \exp \int_{t_0}^{t'} r dt''$$

Note

$$\frac{d}{dt} \left[\left(\exp \int^t r \right) n_+ \right] = r_- \exp \int^t dt' r$$

check:

$$\bullet n_+(t_0) = 1 \cdot n_+(t_0) + 1 \cdot 0$$

$$\begin{aligned}
 \bullet \dot{n}_+(t_0) &= -r n_+(t_0) + \\
 &\quad \exp \left(- \int_{t_0}^{t_0} dt' r \right) \cdot r_-(t_0) \cdot \\
 &\quad \exp \left(+ \int_{t_0}^{t_0} dt'' r \right)
 \end{aligned}$$

$$= -r n_+(t_0) + r_-$$

For $A \ll D$,

$$r_{\pm} = \frac{\sqrt{2}}{2\eta} \cdot \exp\left(-\frac{1}{4D} \mp \frac{A}{D} \cos \Omega t\right)$$

$$\approx r_0 \cdot \left(1 \mp \frac{A}{D} \cos \Omega t\right)$$

Inject $r_1 = \frac{r_0 A}{D}$ in equation for u_+

$$u_+ = \exp -2r_0(t-t_0)$$

$$\left[u(t_0) + \int_{t_0}^t dt' (r_0 + r_1 \cos \Omega t') \exp 2r_0(t'-t_0) \right]$$

$$= \exp -2r_0(t-t_0)$$

$$\left[u(t_0) + \frac{1}{2} (\exp 2r_0(t-t_0) - 1) \right]$$

$$+ \frac{r_1}{\sqrt{\Omega^2 + 4r_0^2}} \cos(\Omega t - \phi) \exp 2r_0(t-t_0) - \cos(\Omega t_0 - \phi) \left. \right]$$

$$\tan \phi = \frac{\Omega}{2r_0}$$

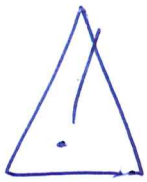
Next: auto correlation function
 \longrightarrow
 power spectrum.

$$\langle X(t) \rangle = x_+ u_+(t) + x_- u_-(t) \\ = u_+ - u_- \quad \text{for } x_{\pm} = \pm 1$$

$$\langle X(t) \cdot X(t+T) | x_0, t_0 \rangle = \\ = \sum_{s, s'} s s' P(s, t+T | s', t) P(s', t | x_0, t_0) \quad \text{Markov property} \\ = u_+(t+T | +1, t) \cdot u_+(t | x_0, t_0) \\ - u_+(t+T | -1, t) \cdot u_-(t | x_0, t_0) \\ - u_-(t+T | +1, t) \cdot u_+(t | x_0, t_0) \\ + u_-(t+T | -1, t) \cdot u_-(t | x_0, t_0).$$

- use $u_- = 1 - u_+$
- substitute above result
- limit $t_0 = -\infty$

$$\langle X(t) X(t+T) | x_0, t_0 \rangle = \\ \exp(-2r_0 T) + 2r_1 + \\ 4r_1^2 [-\exp(-2r_0 T) \cos^2(\Omega t - \varphi) \\ + \cos(\Omega(t+T) - \varphi) \cdot \cos(\Omega t - \varphi)].$$



explicit time-dependence,
even for $t_0 \rightarrow -\infty$.

Reason: periodic driving
defines a phase $\varphi = \Omega t$

• average over phase φ .

$$C(T) = \langle x(t) x(t+T) \rangle_{x_0, t_0, \varphi} =$$

$$\underbrace{\exp(-2r_0 |\tau|) [1 - 2r_1^2]}_{\text{decaying component}} + \underbrace{2r_1^2 \cos \Omega T}_{\text{periodic component}}$$

• Wiener-Khinchin theorem.

$$S(\omega) = C(T) \sim \text{exists.}$$

\equiv power spectral density $S(\omega)$ $\equiv \langle |x(\omega)|^2 \rangle$

A physicist's proof

$$\tilde{x}(\omega) = \int_{-\infty}^{\infty} dt x(t) \exp(-i\omega t)$$
$$x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega \tilde{x}(\omega) \exp(i\omega t)$$

$$C(T) = \langle x(t) \cdot x(t+T) \rangle_t = \left(\frac{1}{2\pi}\right)^2 \int_{-\infty}^{\infty} d\omega_1 \int_{-\infty}^{\infty} d\omega_2$$
$$\langle \tilde{x}(\omega_1) \tilde{x}(\omega_2) \rangle \langle \exp(i(\omega_1 + \omega_2)t) \rangle_t \cdot \exp(i\omega_2 T)$$
$$= \left(\frac{1}{2\pi}\right)^2 \int_{-\infty}^{\infty} d\omega_1 \langle \tilde{x}(\omega_1) \tilde{x}(-\omega_1) \rangle \cdot \exp(-i\omega_1 T)$$

Example: Spectral power density
of Ornstein-Uhlenbeck process

$$T\dot{x} = -x + \xi$$

$$(1 + Tj\omega)\tilde{x}(\omega) = \tilde{\xi}(\omega)$$

$$S_x(\omega) = \langle |\tilde{x}(\omega)|^2 \rangle$$

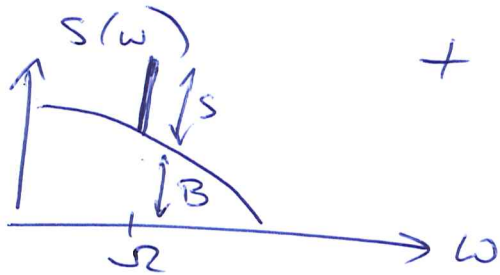
$$= \left| \frac{1}{1 + Tj\omega} \right|^2 \langle |\tilde{\xi}(\omega)|^2 \rangle$$

$$= \frac{2D}{(\omega)^2 + 1}$$

\equiv Lorentzian

Power spectrum:

$$S(\omega) = (1 - 2r_1^2) \underbrace{\frac{4r_0}{4r_0^2 + \omega^2}}_{\text{Lorentzian shape}} + \frac{2\pi r_1^2}{4r_0^2 + \omega^2} (S(\omega - \Omega) + S(\omega + \Omega))$$



• $SNR = \frac{S}{B}$

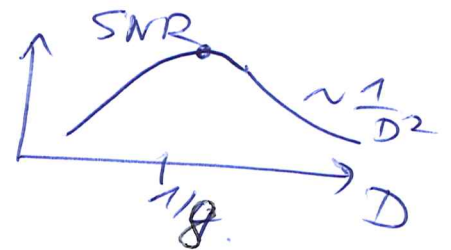
Double-check:
Which limit is taken here?

$$= \pi \frac{r_1^2}{r_0}$$

$$= \pi \cdot \left(\frac{r_0 A}{D}\right)^2 \frac{1}{r_0}$$

$$r_0 = \frac{\sqrt{2}}{2\epsilon} \exp\left(-\frac{1}{4D}\right)$$

$$= \frac{1}{\sqrt{2}} \frac{A^2}{D^2} \exp\left(-\frac{1}{4D}\right)$$



• Spectral amplification factor

$$SPA = \frac{A_1(D)^2}{A^2}$$

$$2x(t) = A_1(D) \cos(\Omega t - \varphi)$$

$$A_1(D) = \frac{A}{D} \frac{2r_0}{\sqrt{4r_0^2 + \Omega^2}}$$

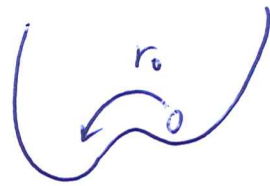
$$= \frac{4r_0^2}{D^2 (4r_0^2 + \Omega^2)}$$

\Rightarrow SPA attains maximum
at $\Omega = 2r_0$.

\Rightarrow Intuitive interpretation



$$\frac{1}{r_0} \approx \frac{T}{2}$$



$$\frac{1}{r_0} \approx \frac{T}{2}$$

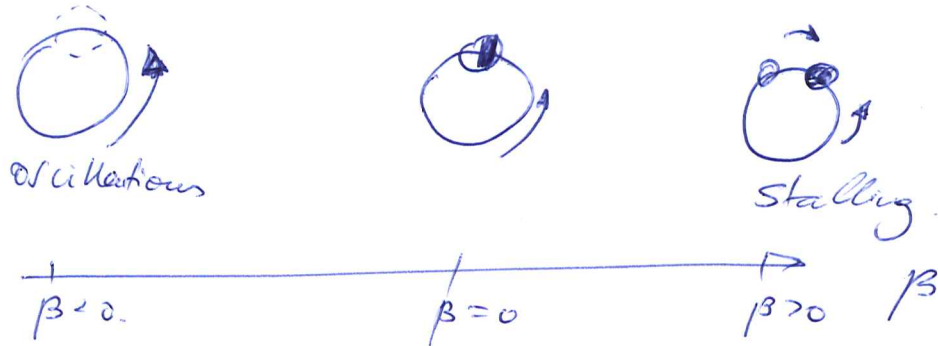
Biological exple

Mechano-receptor in crayfish

- predator = periodic pressure signal
turbulence = noise.
- sensory detection optimal
at physiological noise
levels.
- experimental test by
playing signal + variable
noise.
and recording neuron spiking

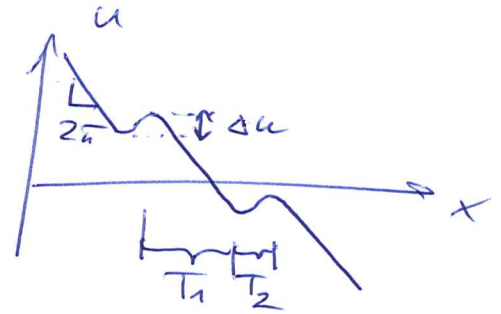
Coherence resonance

Recall: saddle-node bifurcation



Now with noise:

$$\dot{X} = -\frac{\partial u}{\partial X} + \xi$$



period: $T = T_1 + T_2 \equiv$ random variable

• T_1 : $\dot{X} = 2u + \xi$, $\langle \xi(t)\xi(t') \rangle = 2D \delta(t-t')$
 $\langle T_1 \rangle \approx 1$, $\langle T_1^2 \rangle - \langle T_1 \rangle^2 \approx 2D$.

• T_2 : Poisson point process:
 exponential distribution
 of waiting times

$$\langle T_2 \rangle = \frac{1}{r}, \quad \langle T_2^2 \rangle - \langle T_2 \rangle^2 = \frac{1}{r^2}$$

$$r = \frac{1}{a} \exp\left(-\frac{b}{D}\right)$$

• $\frac{\langle T^2 \rangle}{\langle T \rangle^2} = 2D + a \exp\left(+\frac{b}{D}\right)$

