

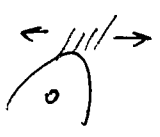
4

Nonlinear dynamics of active oscillators

Examples of active, biological oscillators.

① circadian rhythm $T = 24 \text{ h}$

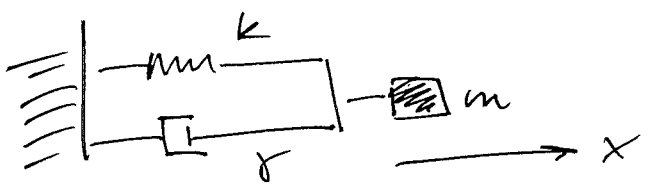
② single heart muscle cell $T = 1 \text{ s}$

③ hair bundle in the inner ear
 $T = 0.1 \text{ ms}$

→ active oscillators that oscillate spontaneously.

Today: Minimal model for active oscillators.

Example of passive oscillator:
harmonic oscillator



m ... mass
 k ... spring constant
 γ ... friction coefficient

force balance:

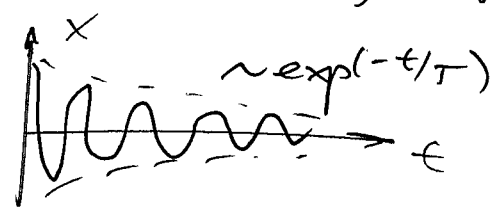
$$m\ddot{x} + \gamma\dot{x} + kx = 0$$

inertia friction elasticity

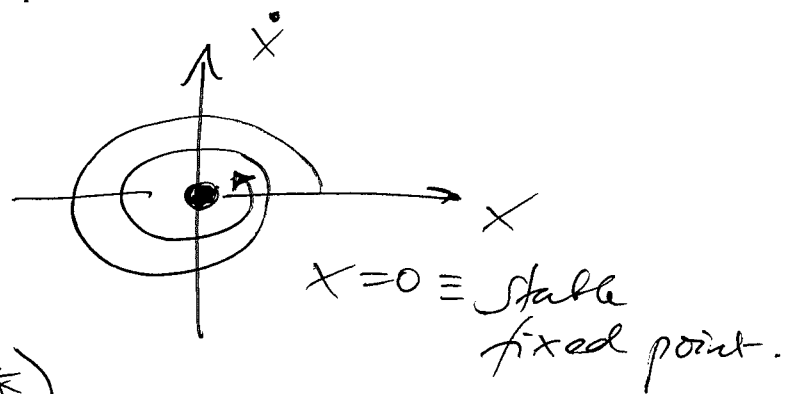
• Solutions are damped oscillations

(*) $x(t) = x(0) \exp(-t/T) \cos \omega_0 t$
 with

$T = m/2\gamma$... damping time
 $\omega_0 \approx \sqrt{k/m}$... resonance freq.
 for $\gamma^2 < 4cm \cdot k$



Phase dynamics.



Exercise: check (*).

Thought experiment:

Q: What happens for $\gamma < 0$?

A: Oscillation amplitude grows as $\exp(+t/T)$.

This requires energy input.

Active oscillators can be described by nonlinear friction $\gamma = \gamma(x)$ that can become negative.

Example: Van-der-Pol oscillator.

Choice $\gamma = \mu(1 - 2x^2)$ yields

$$(**) \quad m\ddot{x} + \mu(1 - 2x^2)\dot{x} + kx = 0.$$

Ansatz:

$$x(t) \approx a \cos \omega_0 t.$$

What sign has γ ?

Trick: average over one oscillation cycle $\Rightarrow \langle \gamma \rangle = \mu(1 - a^2).$

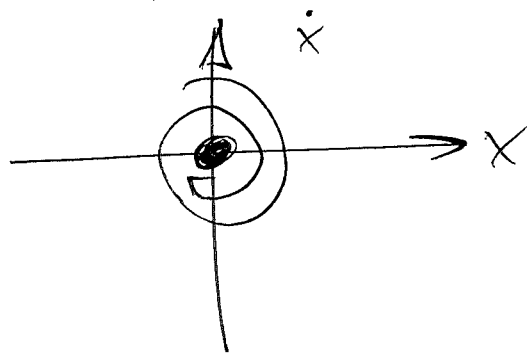
$\Rightarrow \langle \gamma \rangle < 0$ for $a < 1^{1/2}$
 \Rightarrow amplitude grows

$\langle \gamma \rangle > 0$ for $a > 1^{1/2}$
 \Rightarrow amplitude decays

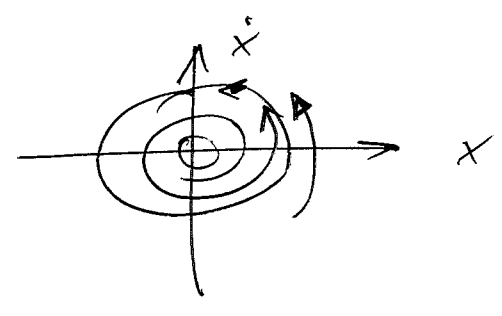
\Rightarrow stable limit cycle oscillations
 with $a = 1^{1/2}$, $\omega_0 \approx \sqrt{k/m}$
 if $\mu > 0$.

Phase dynamics

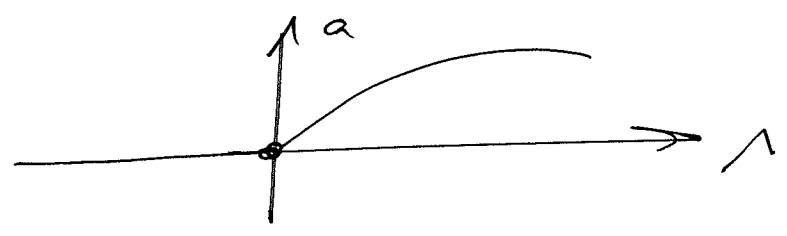
$\Lambda < 0$
(passive)



$\Lambda > 0$
(active)



The onset of spontaneous oscillations for $\Lambda \geq 0$ corresponds to a Hopf bifurcation with control parameter Λ .



This transition from a stable fixed point $a=0$ for $\Lambda < 0$ to spontaneous oscillations for $\Lambda > 0$ is called a Hopf bifurcation.

Normal form of Hopf bifurcation

$$(*) \quad \dot{Z} = i(\omega_c - \omega_1 |Z|^2)Z + \mu(1 - |Z|^2)Z.$$

with complex oscillator variable

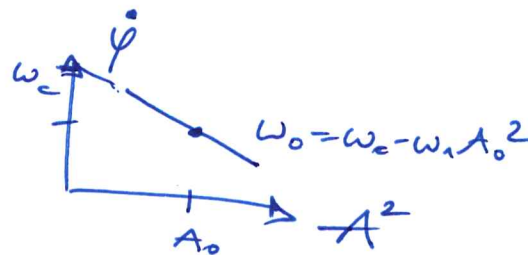
$$Z = A \exp i\varphi.$$

We can rewrite (*) as

$$\dot{A} = \mu(1 - A^2)A.$$

$$\approx \mu A_0^2(A_0 - A) \quad \text{if } |A - A_0| \ll A_0 = 1$$

$$\dot{\varphi} = \omega_c - \omega_1 A^2$$



$\omega_1 \equiv$ non-Isodromy.

Comment:

How to construct Z ?

For periodic signal $x(t) = A \cos \omega_0 t$.

Set $Z = x(t) + i \dot{x}(t) / \omega_0$.

Q: What happens at $\lambda=0$?

A: The oscillator is highly sensitive to external driving at the resonance frequency

$$m\ddot{x} + \mu(0 - 2\mu x^2)\dot{x} + kx = f \cos \omega_0 t. \quad \text{ext. driving.}$$

Ansatz: $x(t) = a \sin \omega_0 t.$

$$-2\mu a^2 \underbrace{\sin^2 \omega_0 t}_{\Rightarrow \text{averages } 1/2} - \omega_0 a \cos \omega_0 t = f \cos \omega_0 t$$

$$\mu \omega_0 a^3 = f.$$

$$a \sim f^{1/3}$$

$$\text{gain} = \frac{a}{f} \sim f^{-2/3}.$$

Application:

The dynamics of hair bundles in our inner ear can be

mapped onto (**).

Hair bundles are tuned at $\lambda=0$ to amplify faint sounds (low f) more than loud sounds (high f).

Poincaré - Bendixon - Theorem.

4-8

Existence of closed orbits in 2D.

$$(\dot{x}, \dot{y}) = f(x, y), \quad f \in C^1$$

If $R \subseteq \mathbb{R}^2$ is a trapping region

- closed, bounded

- No trajectory can leave R

then R contains either

a fixed point or a
closed orbit.

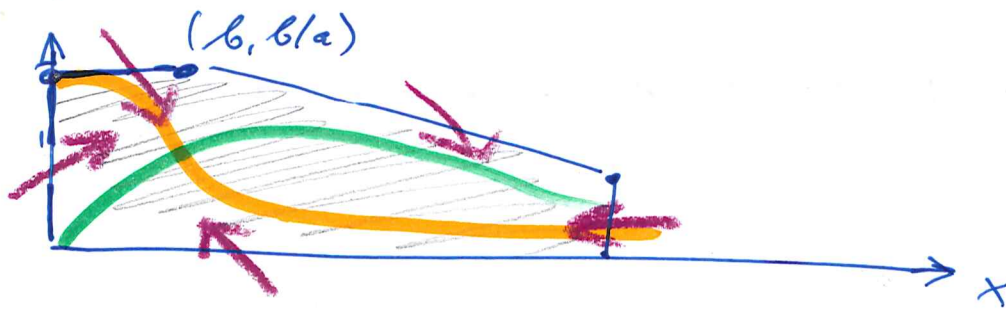
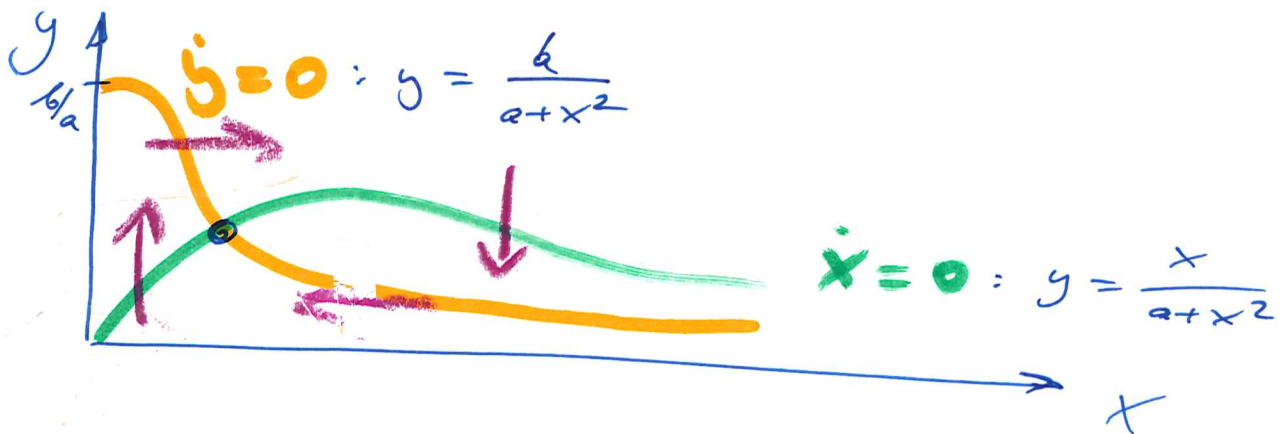
Example:

Glycolysis Model (Selkov, 1968)

$$x \equiv [\text{ADP}]$$

$$y \equiv [\text{Fructose-6P}]$$

$$\begin{aligned} \dot{x} &= -x + ay + x^2y \\ \dot{y} &= b - ay - x^2y \end{aligned}$$



fixed point

$$x^* = b, \quad y^* = \frac{b}{a+b^2}$$

$$J = \begin{pmatrix} 1+2x^*y^* & a+x^{*2} \\ -2x^*y^* & -(a+x^{*2}) \end{pmatrix}$$

$$\Rightarrow \Delta = a+b^2 > 0$$

$$T = \frac{-b^4 + (2a-1)b^2 + (a+a^2)}{a+b^2}$$

$T > 0 \Rightarrow$ unstable
 $T < 0 \Rightarrow$ stable

